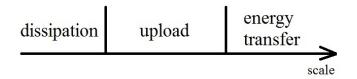
Checking the stability of solutions in shell models using symbolic computation

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Turbulence representation



Physical space - processes are mixed, present at every point

Space of scales (wave vectors) - processes are separated

Shell turbulence models

The wavenumber axis is split into progressively expanding octaves

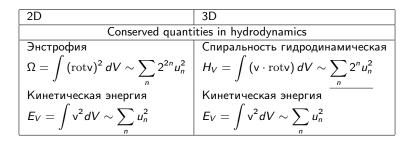
$$k_n < |\mathbf{k}| < k_{n+1}, k_n = 2^n k_0$$

General view of shell equations

$$\frac{d_t X_n}{dt} = \sum_{ij} Q_{nij} X_i X_j - K_n X_n + f_n \tag{1}$$

 Q_{nij} – nonlinear interaction matrix $K_n X_n \sim k_n^2 X_n$ – dissipative term f_n – external forces

Table 1. Stored quantities



Hydrodynamic system equations

GD turbulence equations in the Boussinesq approximation in dimensionless form

$$\frac{\partial \mathsf{v}}{\partial t} + (\mathsf{v} \bigtriangledown) \,\mathsf{v} = - \bigtriangledown \, \mathsf{p} + \mathsf{R} \mathsf{e}^{-1} \,\triangle \,\mathsf{v},\tag{2}$$

 Re – Reynolds number

Conserved quantities

Kinetic energy
$$E = \int (v^2) dV$$

Enstrophy $\Omega = \int (\text{rotv})^2 dV$
Hydrodynamic helicity $H_V = \int (v \cdot \text{rotv}) dV$

Shell model of hydrodynamic turbulence

Turbulent convection shell model including nonlocal interactions

$$\frac{d_t u_n}{dt} = \sum_{ij} S_{nij} u_i u_j - Re^{-1} k_n^2 u_n, \tag{3}$$

Conservation laws

Kinetic energy

$$\frac{dE}{dt} = 0 = \frac{d}{dt} \sum_{n} u_{n}^{2} = 2 \sum_{n} u_{n} \frac{du_{n}}{dt}$$

Enstrophy

$$\frac{d\Omega}{dt} = 0 = \frac{d}{dt} \sum_{n} 2^{2n} u_n^2 = 2 \sum_{n} 2^{2n} u_n \frac{du_n}{dt}$$

Hydrodynamic helicity

$$\frac{dH_V}{dt} = 0 = \frac{d}{dt} \sum_n 2^n (-1)^n u_n^2 = 2 \sum_n 2^n (-1)^n u_n \frac{du_n}{dt}$$

The resulting equations

The system consists of three families of equations for s[i, j]

$$(s_{i,j} + s_{j,i}) + 2^{i} (s_{-i,j-i} + s_{j-i,-i}) + 2^{j} (s_{i-j,-j} + s_{-j,i-j}) = 0,$$

$$(s_{i,j} + s_{j,i}) + 2^{3i} (s_{-i,j-i} + s_{j-i,-i}) + 2^{3j} (s_{i-j,-j} + s_{-j,i-j}) = 0,$$

$$(s_{i,j} + s_{j,i}) + (-1)^{i} 2^{2i} (s_{-i,j-i} + s_{j-i,-i}) + (-1)^{j} 2^{2j} (s_{i-j,-j} + s_{-j,i-j}) = 0$$

(4)

The system is also complemented by symmetries of the following form

$$s_{i,j}-s_{j,i}=0 \tag{5}$$

Scheme for calculating equations in a symbolic package

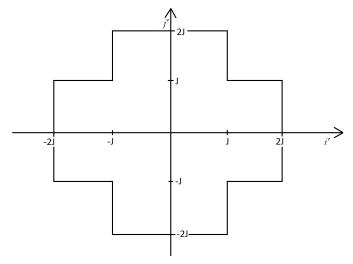


Рис. 1: Region of existence of equations.

Requirements for the existence of stationary solutions with a power-law dependence on the spatial scale

Stationary solutions

$$u = u_0 2^{\frac{p}{q}},\tag{6}$$

Additional equations of the form are added to the system

$$\sum_{i,j} (s[i,j]) 2^{-(i+j)\frac{p}{q}},$$
(7)

where p/q – power law degree.

Interacting shells

Navier-Stokes equation in Fourier space

$$\frac{\partial \widehat{\mathbf{u}}}{\partial t} = i \iint_{\mathbb{R}^{3} \times \mathbb{R}^{3}} S(\mathbf{k}, \mathbf{s}, \mathbf{q}) \,\delta(\mathbf{s} + \mathbf{q} + \mathbf{k}) \,\widehat{\mathbf{u}}(\mathbf{s}, t) \,\widehat{\mathbf{u}}(\mathbf{q}, t) \,d\mathbf{s}d\mathbf{q} - \nu \mathbf{k}^{2} \,\widehat{\mathbf{u}}(\mathbf{k}, t) + \widehat{\mathbf{f}}(\mathbf{k}, t)$$
(8)

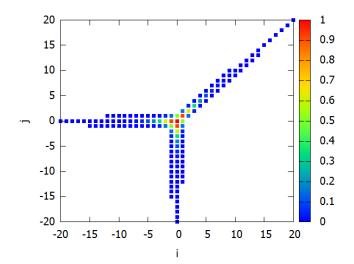
Ranges k, s и q

$$\begin{split} & \mathsf{k} \in (2^{n}; 2^{n+i}) \sim n, \\ & \mathsf{s} \in (2^{n+i}; 2^{n+i+1}) \sim n+i, \\ & \mathsf{q} \in (2^{n+j}; 2^{n+j+1}) \sim n+j, \end{split}$$

Triangle inequality

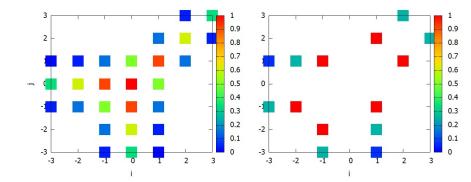
$$2^{n} + 2^{n+i} \ge 2^{n+j}, 2^{n+i} + 2^{n+j} \ge 2^{n}, 2^{n} + 2^{n+j} \ge 2^{n+i},$$
(10)

The probabilities of possible interactions of waves from the nth, (n+i)th, (n+j)th

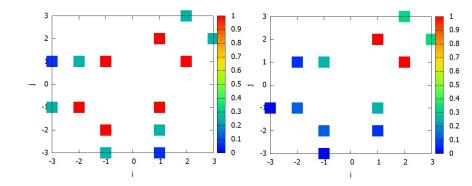


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Interaction probabilities subject to constraints for $\mathsf{J}=\mathsf{3}$



Comparison with the intensities of interactions



Выводы

• The previously developed technology for constructing cascade models of computer algebra methods was supplemented and provides the presence of stationary solutions with any desired exponents. As a result, entire parametric classes of models are generated in an automated mode.

• The developed technique for identifying specific models in which the coefficients of nonlinear interactions would be maximum in magnitude to the probabilities of wave interactions of the corresponding scales.

• In fixed models, a study of power-law solutions for Lyapunov stability was realized. In the examples considered, these solutions turn out to be unstable.

• Further development of this technology presupposes the optimization of computational procedures, since the coordination of the coefficients with the probabilities is calculated for a very long time. And the agreement was not very good.